

# How to bake puffy planets

## Coupling radius inflation with high eccentricity migration

Jiapeng Gao  
Advisor: Gongjie Li



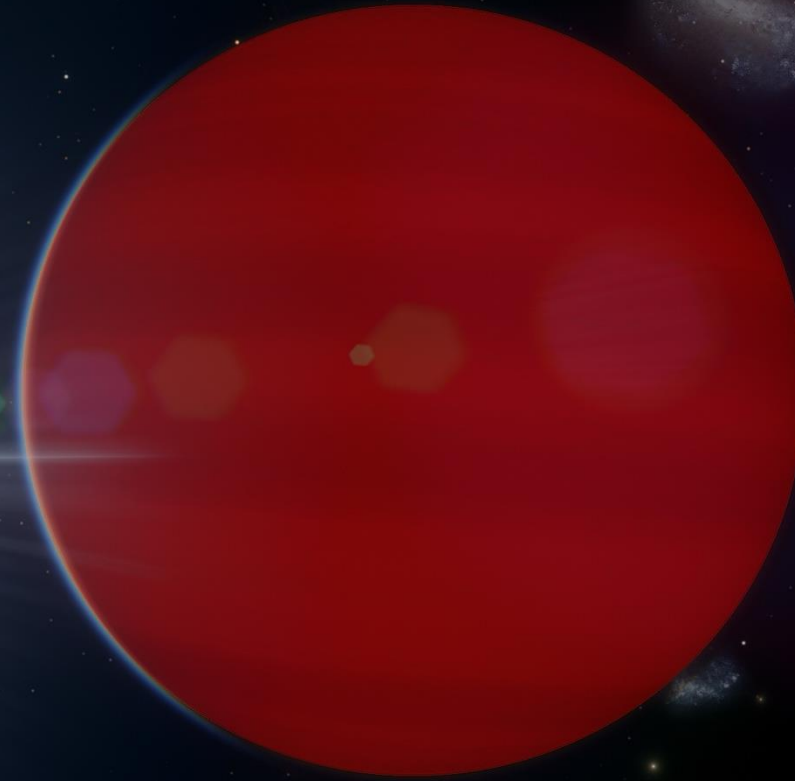
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# Origin of warm puffy planets



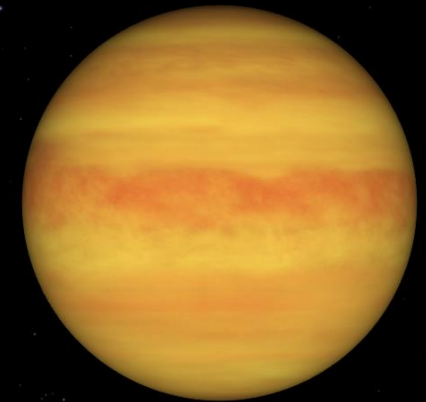
# Origin of warm puffy planets

- WASP-117b -- Saturn mass planet with Jupiter radius
- WASP-117b has only a density of 0.345 g/cm<sup>3</sup>, while Saturn has 0.687 g/cm<sup>3</sup> (Lendl et al. 2014)



EYES ON EXOPLANETS

You are  
515 light-years  
from Earth



WASP-117 b 

A giant planet composed mainly of gas

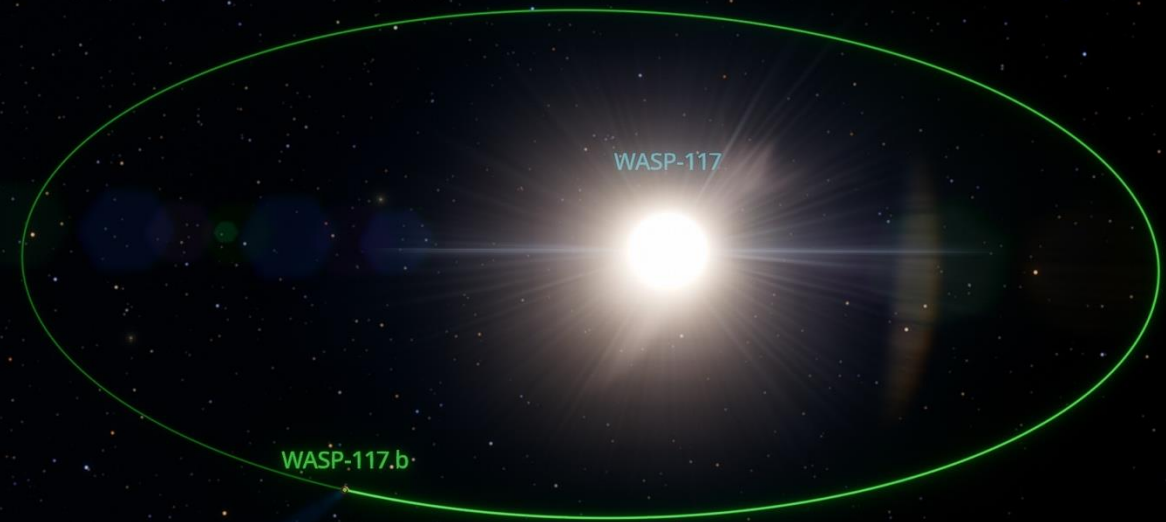
Credit: NASA/Lunar and Planetary Institute

# Puffy planets

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Parameter	Value
Orbital period (days)	10.02165 ± 0.00055
Eccentricity $e$	0.302 ± 0.023
True obliquity $\psi$ (deg)	69.6 <sup>+4.7</sup> <sub>-4.1</sub>
Semi-major axis $a$ (au)	0.09459 <sup>+0.00084</sup> <sub>-0.00079</sub>

Table 1: WASP-117b Orbital Parameters





# Puffy planets

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How does tidal heating affect the structure of the planet, which in turn influence the dynamical evolution?

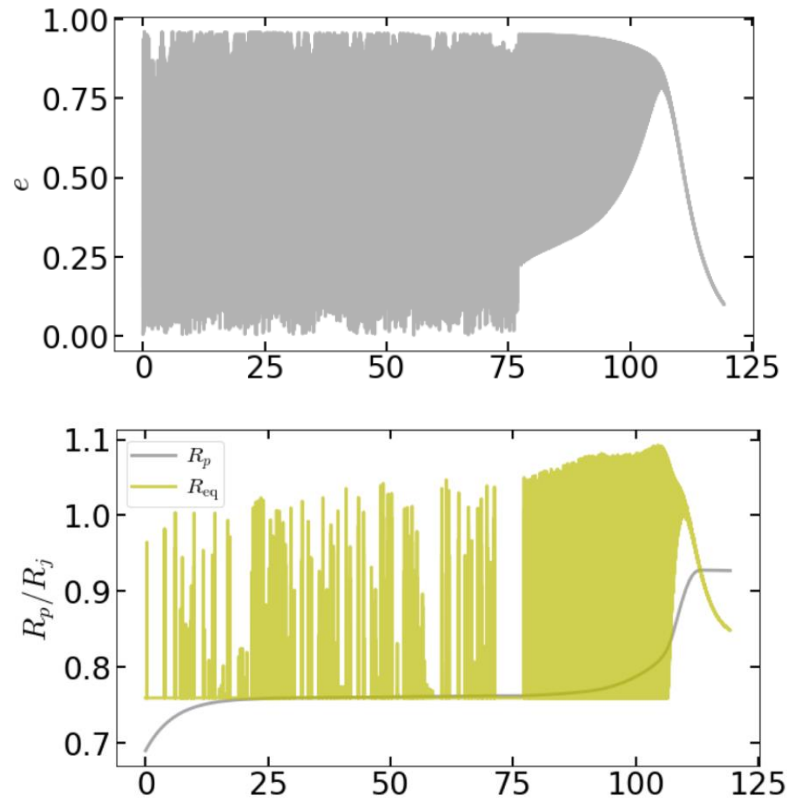
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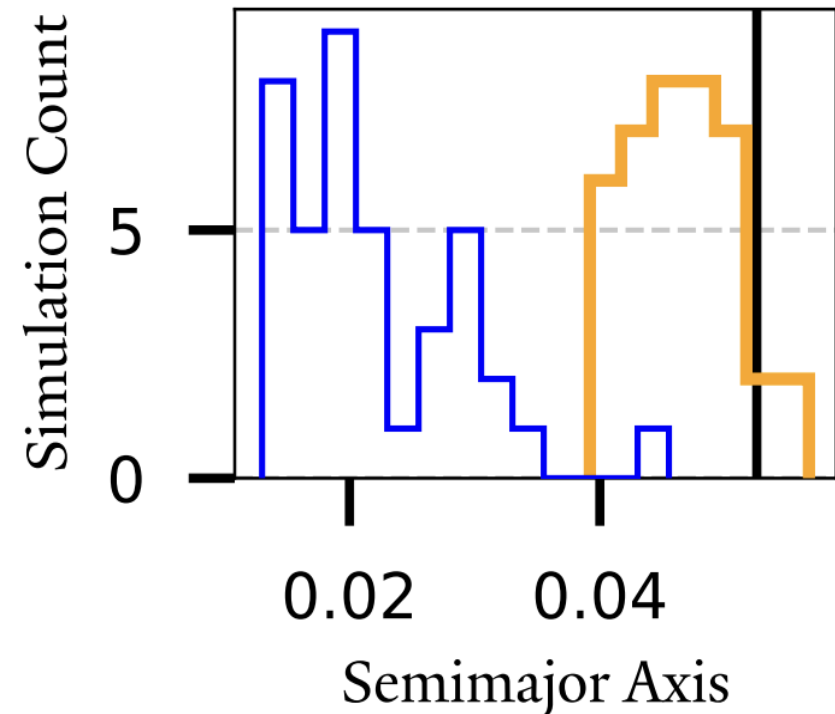
WASP-117.b

# Previous works on coupling radius inflation with tidal migration

- Yu & Dai, 2024:
  - Formation of WASP-107b with high-e migration and radius inflation.



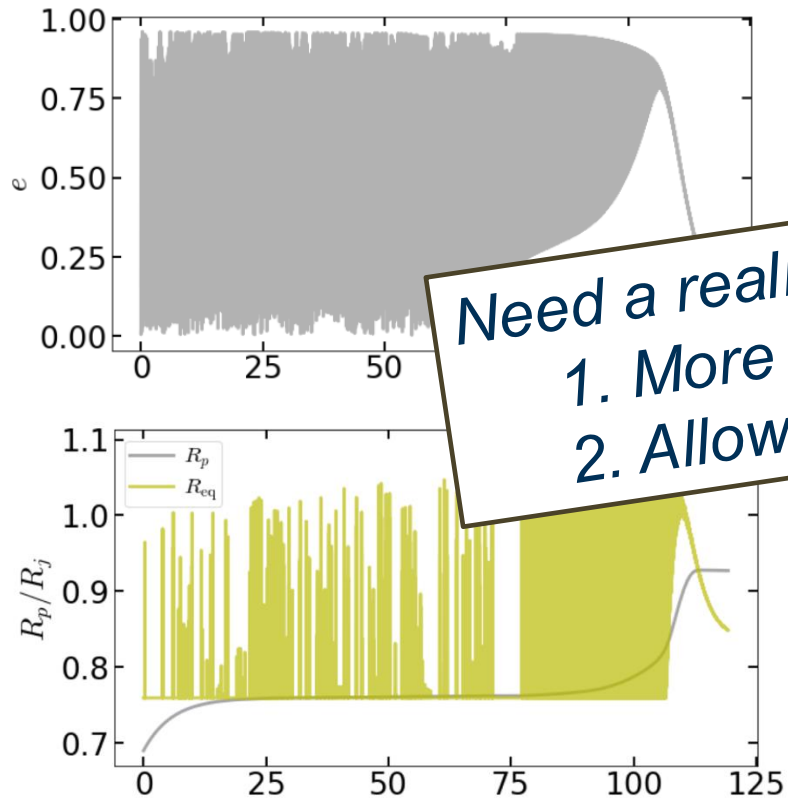
- Lu et al., 2024:
  - Formation of HAT-P-11b with scattering, high-e migration and radius inflation.
  - Radius inflation saved the planet from tidal disruption



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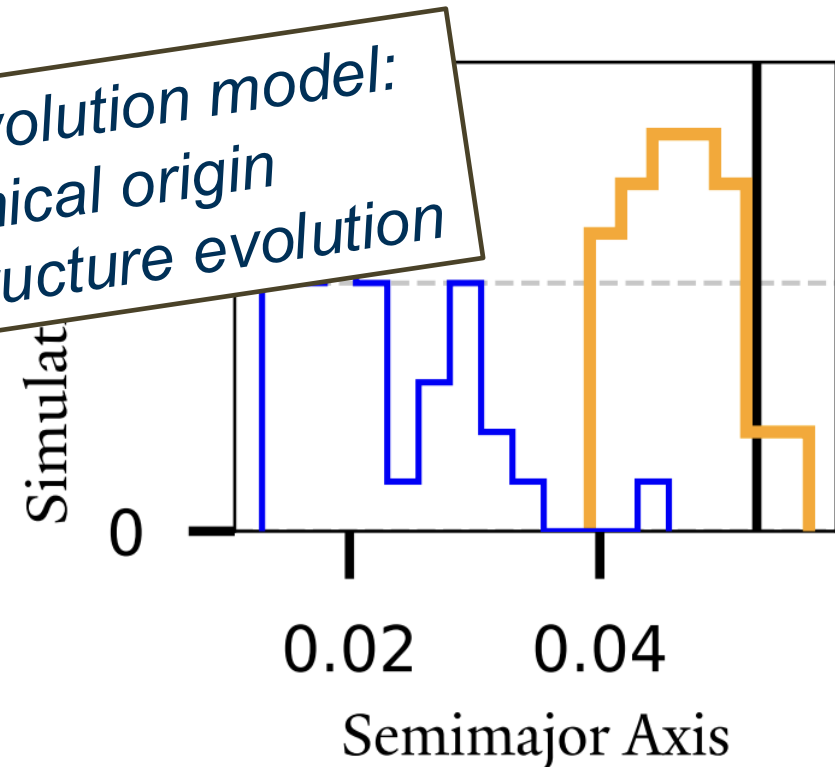
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Need a realistic structure evolution model:  
1. More accurate dynamical origin  
2. Allows a probe for structure evolution

- Lu et al., 2024:

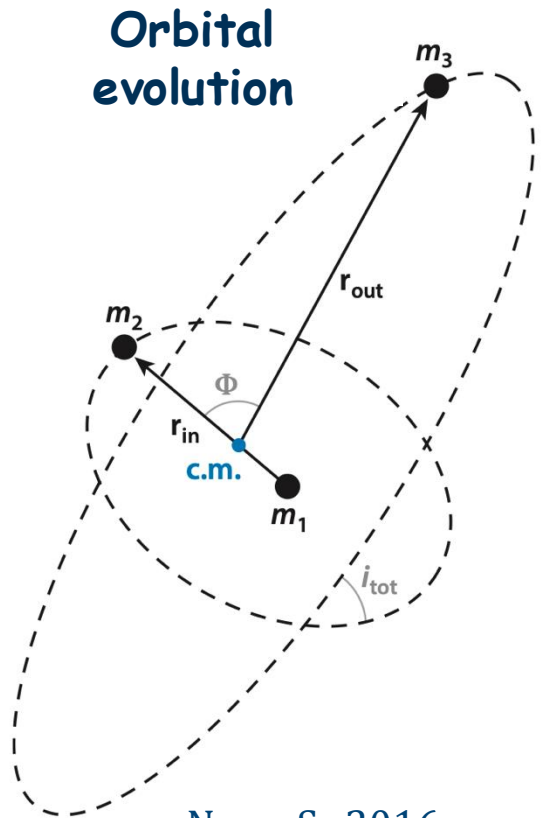
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# How are puffy planets formed?

- First combination of realistic interior structure evolution with high-e migration.



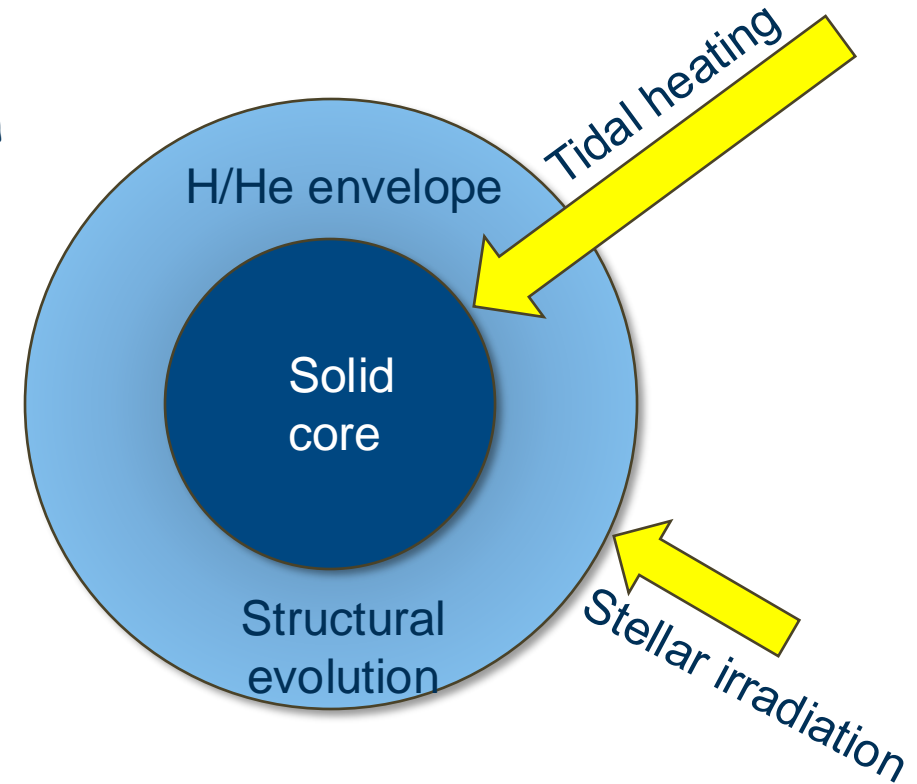
Naoz, S., 2016.

Octupole L-K+ Short range forces

Tidal heating rate & Surface heating rate



Radius inflation

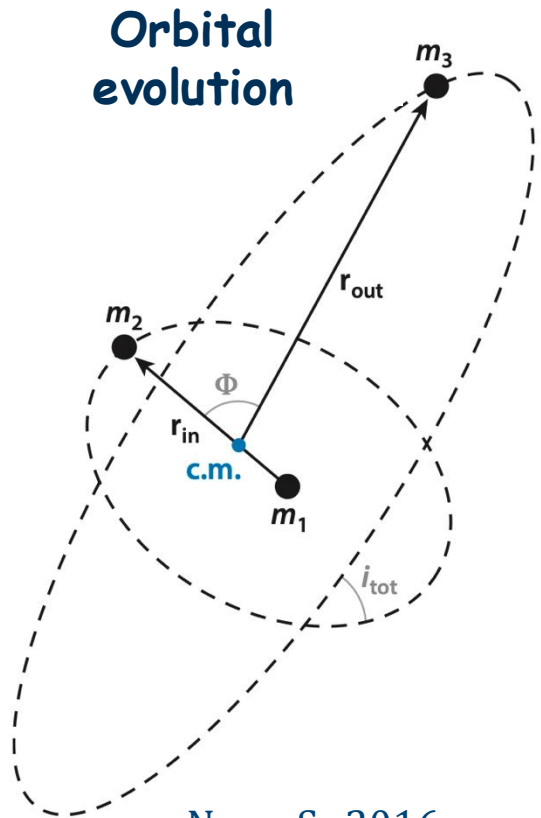


MESA

(Paxton et al., 2011)

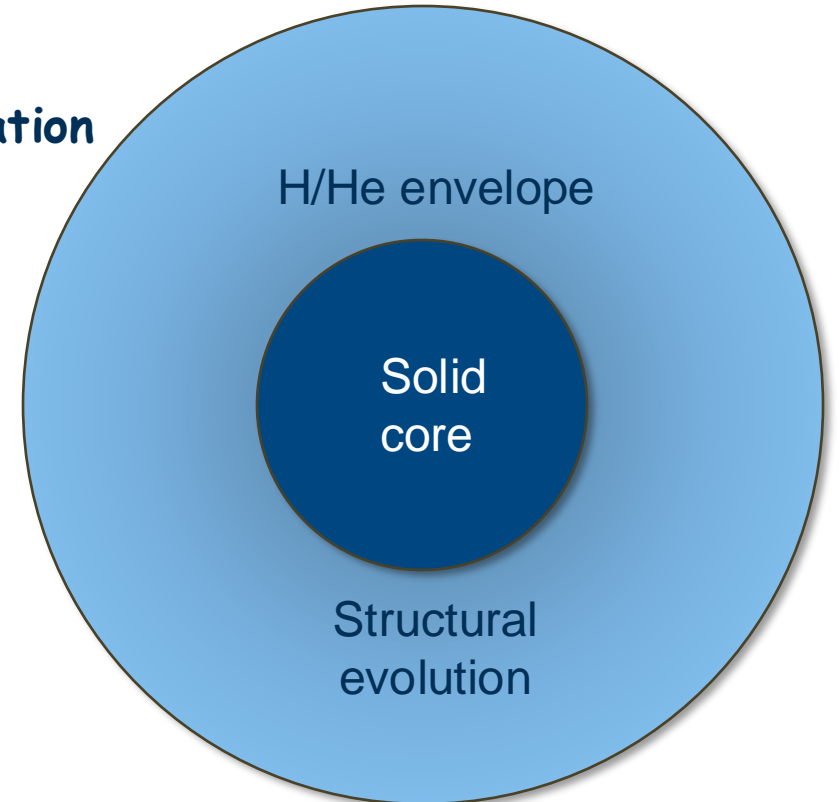
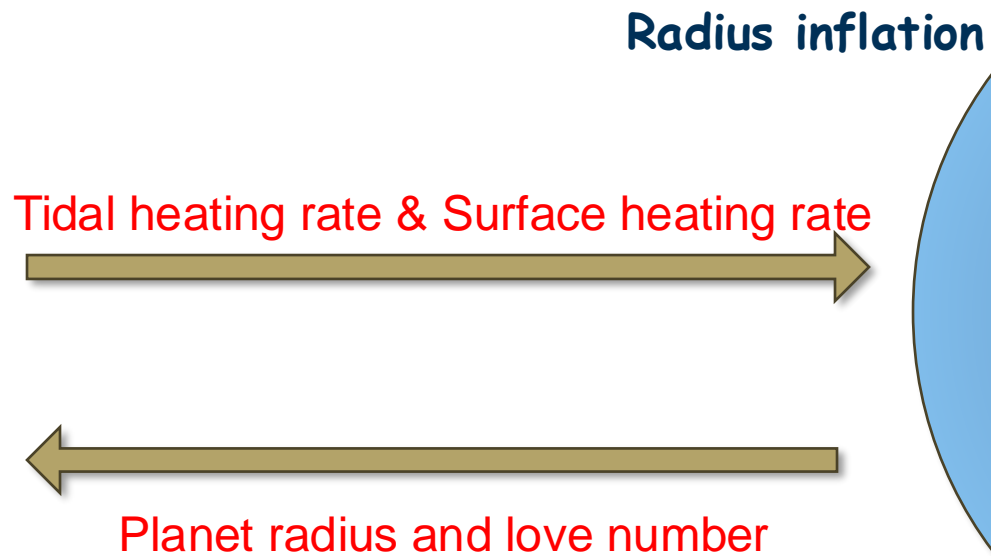
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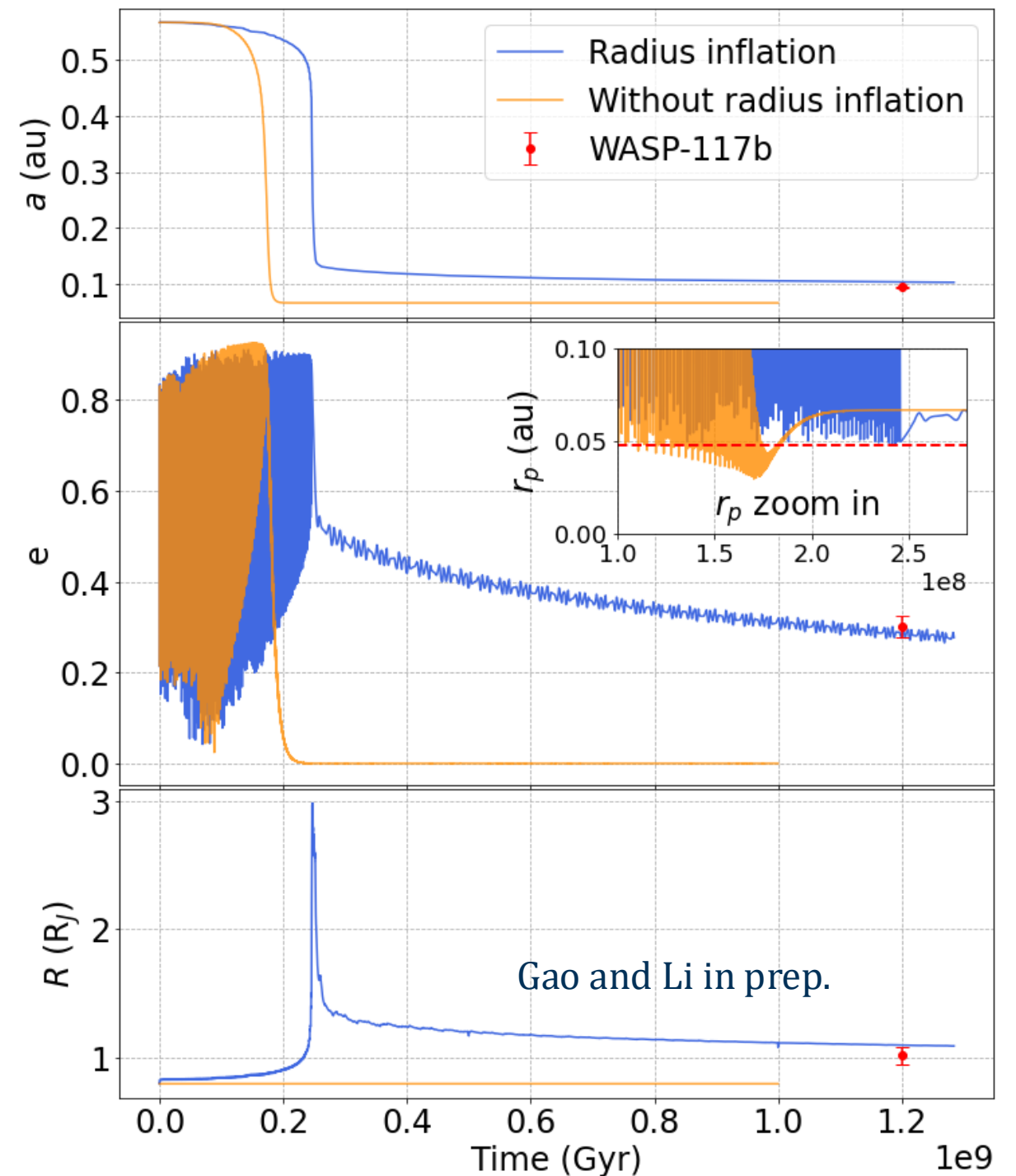


**MESA**  
(Paxton et al., 2011)

# Radius inflation

- Eccentricity excitation increases tidal heating rate
- Tidal heating inflates the planet, triple the radius.
- L-K suppressed, semi-major axis decrease sharply with residual eccentricity
- Planet quickly cools down, orbit slowly circularize.

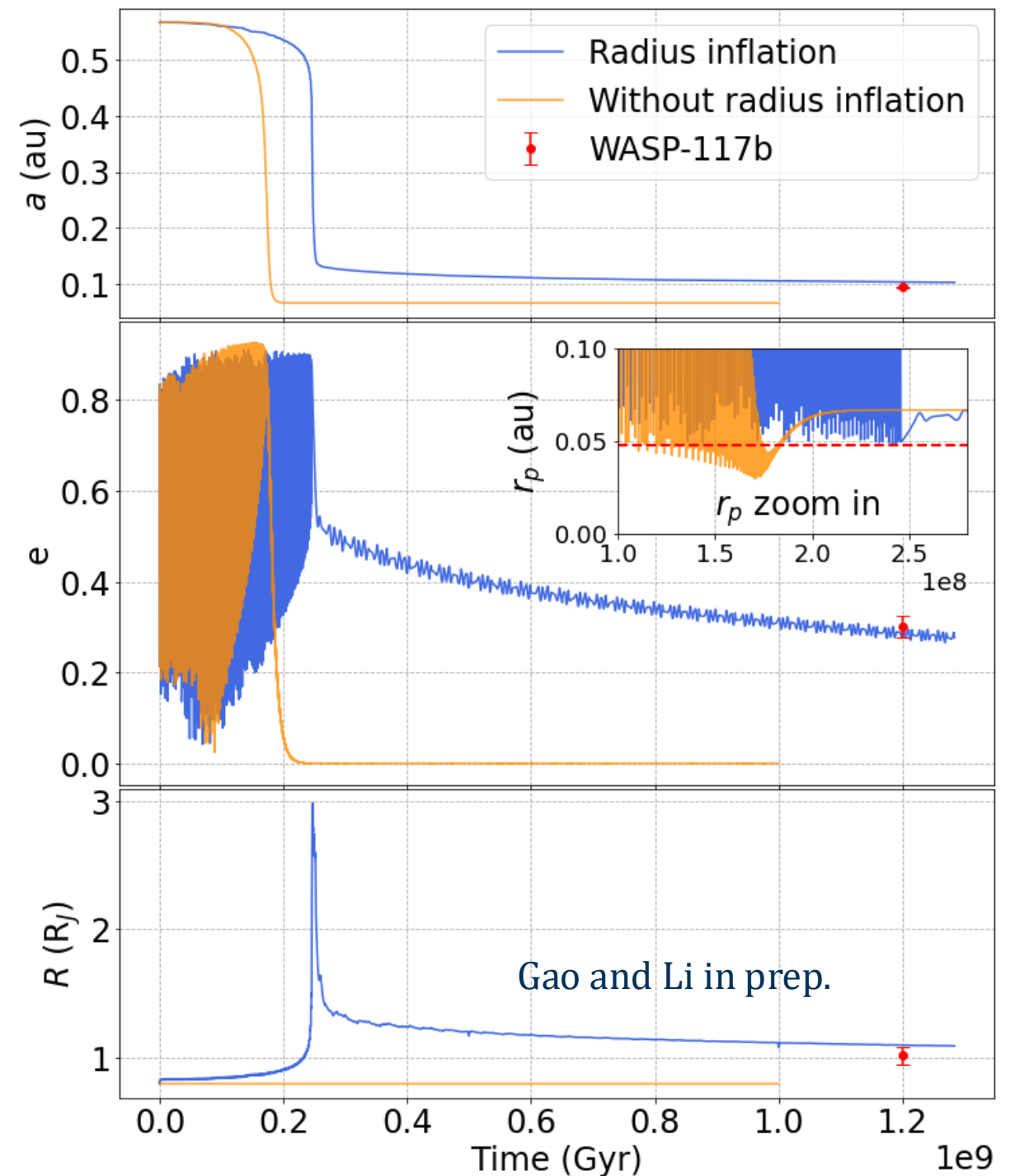
Results agree well with observation 😊



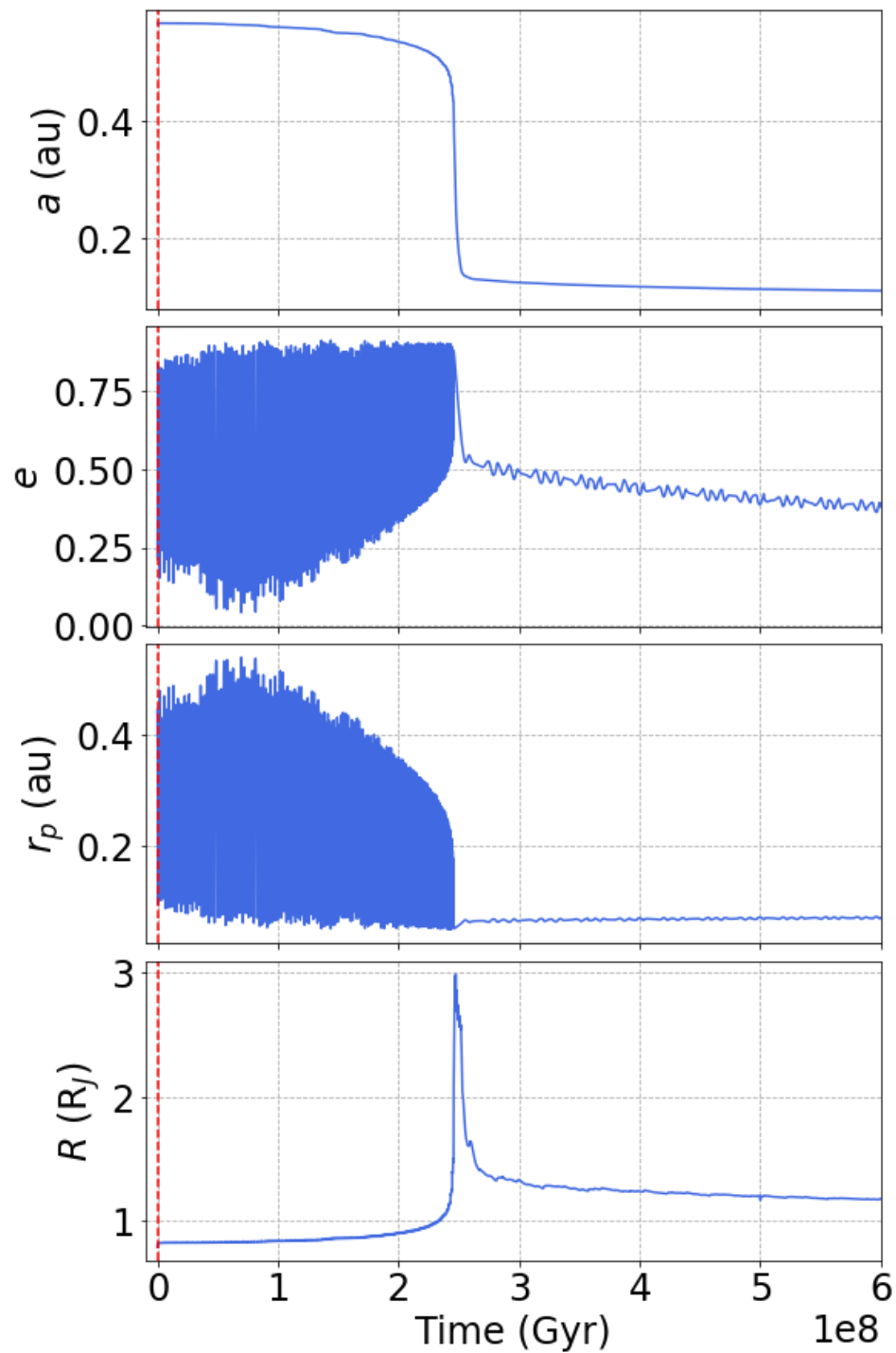
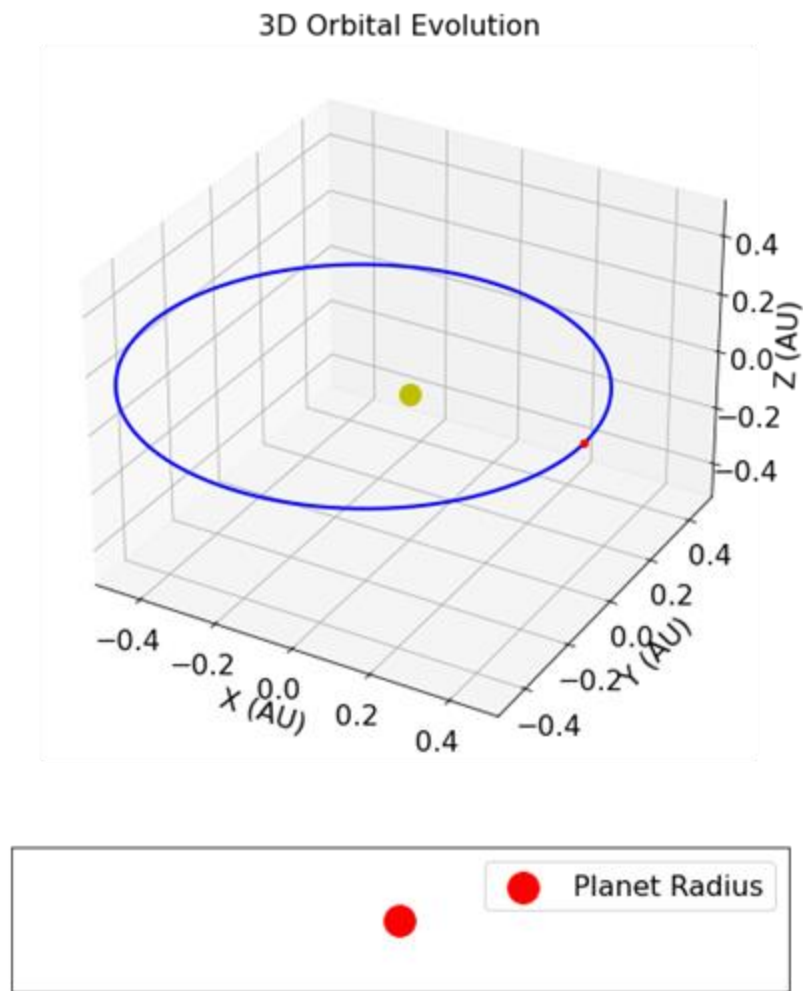
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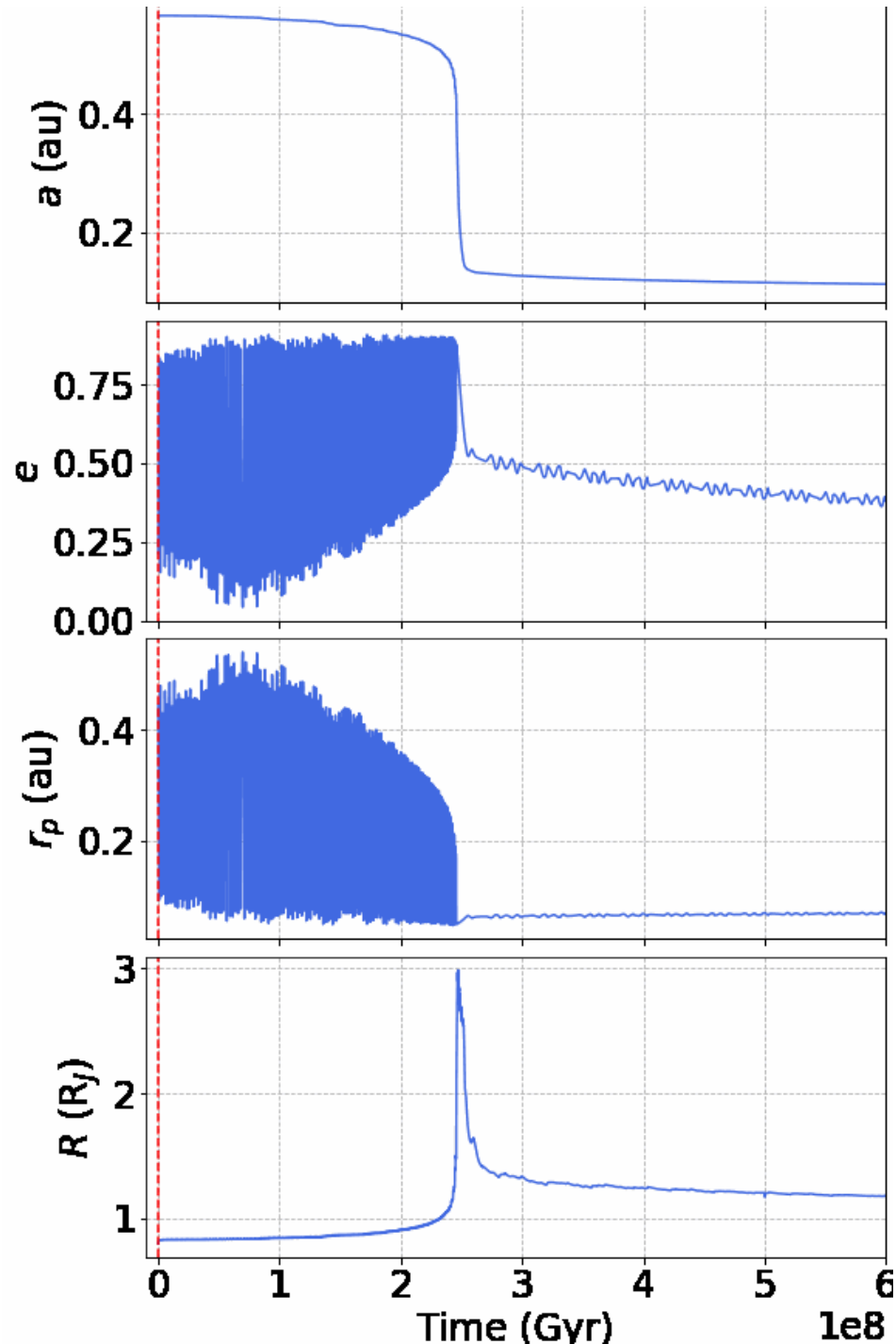
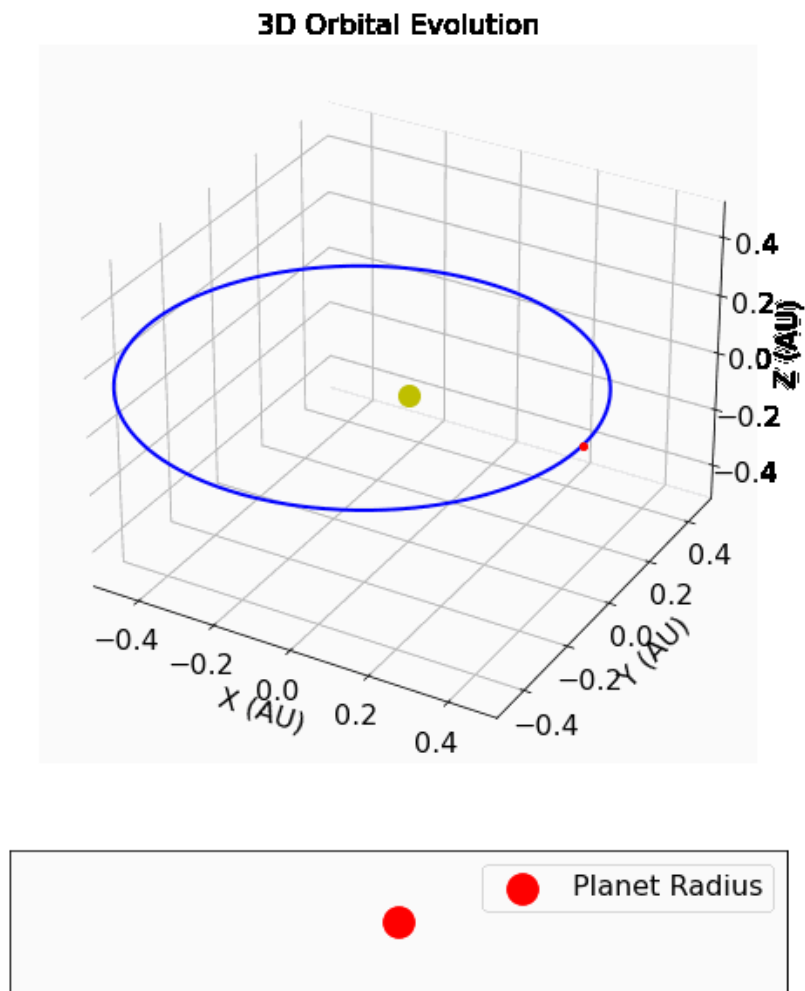
*In short, we will have **larger** final  $a$ ,  $e$  and **radius** with radius inflation.*



# Radius inflation



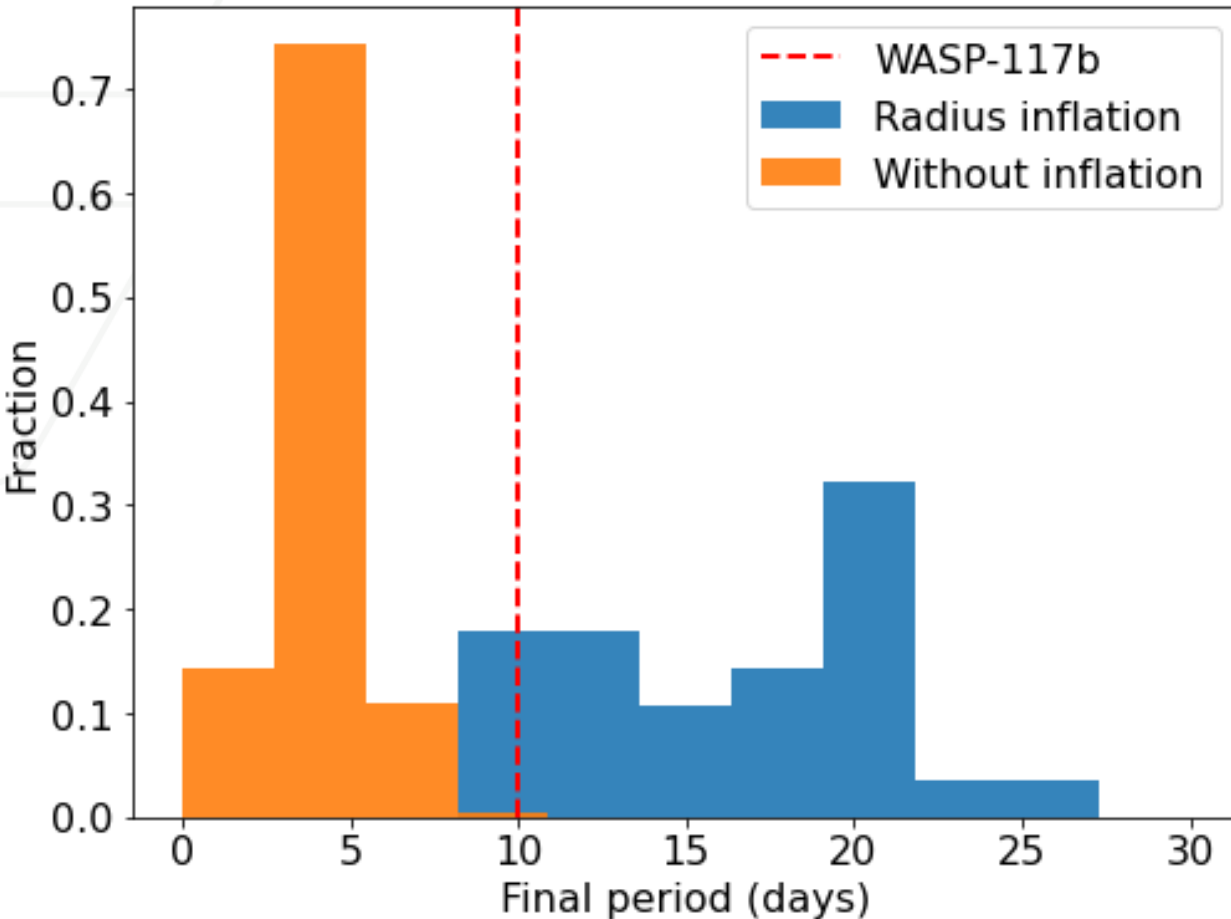
# Radius inflation





# Parameter dependencies

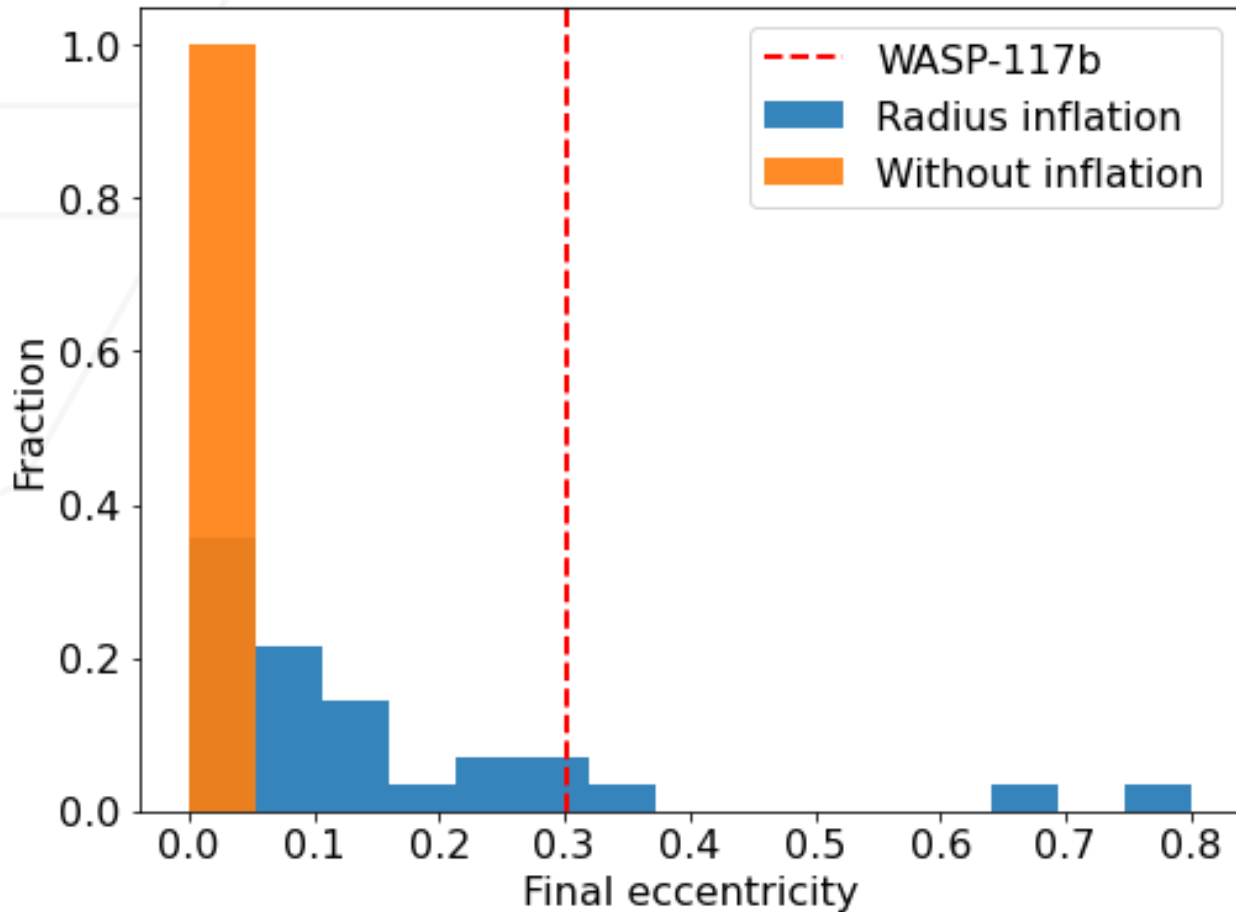
- Final orbital period distribution



- Radius inflation increases the final orbital period to  $\sim 10$  days.
- $\sim 10$  days is exactly what we need for warm Saturn WASP-117b.

# Parameter dependencies

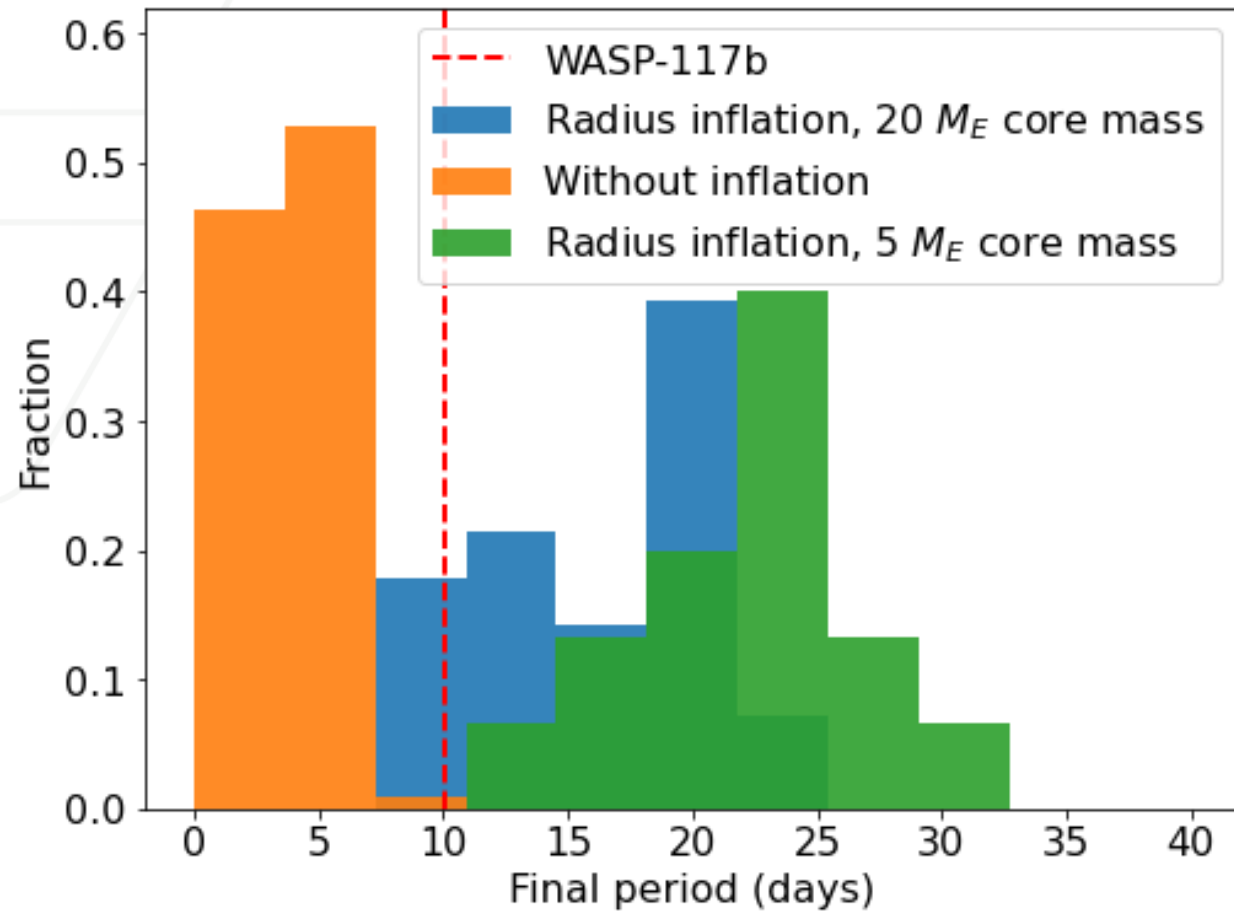
- Final eccentricity distribution



- Radius inflation induces larger final eccentricity, due to larger pericenter distance when Kozai suppressed
- Explains the observed residual eccentricity.

# 5 $M_E$ core mass simulation

- Final orbital period distribution



- Larger final period than 20  $M_E$  core mass case.
- Lower core mass leads to higher radius inflation, thus suppresses L-K effect at even larger pericenter distance.

# Conclusion

- Tidal heating is capable of explaining the radius inflation of **warm** puffy planets.
- Radius inflation suppressed L-K during migration of puffy planets, leading to residual eccentricity and larger orbital periods compared with fixed radius.
- Smaller core mass generates larger final radius, larger final orbital periods.
- More transit (eg. **TESS**) and RV observation for precise determination of the radii and masses of inner planets.
- Long term RV measurements or astrometric measurements, (eg. **Gaia**) for characterizing the outer planets.
- Atmosphere chemical abundance observation (eg. **JWST**) also helps us determine the interior structure and core mass. (Sing et al., 2024)



**Thank you**

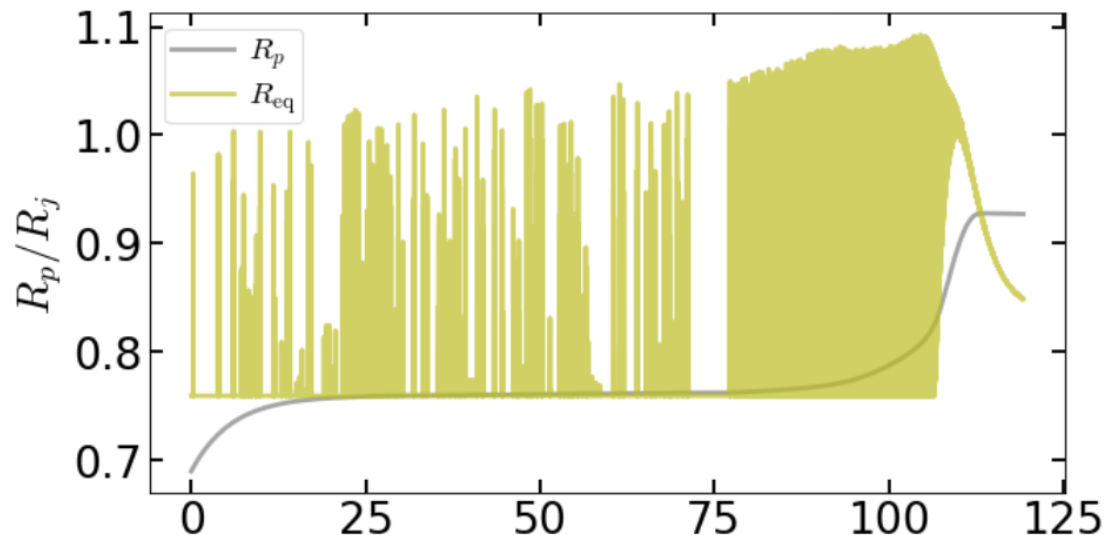
# Caveats of previous works

- Yu & Dai, 2024:
  - The radius inflation model is an analytical evolution between equilibrium radii

$$\frac{dR}{dt} = \frac{R_{\text{eq}} - R}{\tau},$$

$$\tau = \begin{cases} \tau_{\text{def}}, & R > R_{\text{eq}}, \\ \tau_{\text{inf}}, & R < R_{\text{eq}}. \end{cases}$$

- Constant timescale should depend on tidal inflation.

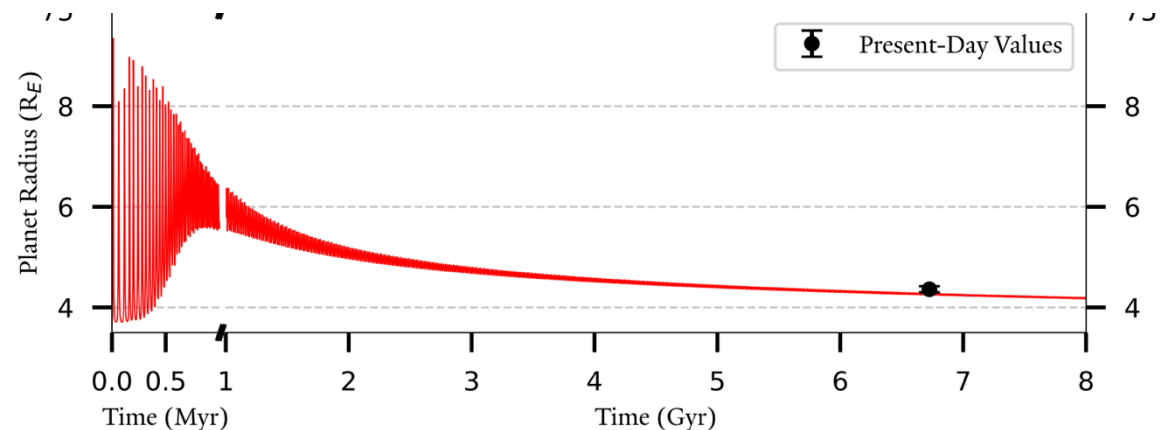


- Lu et al., 2024:
  - Used a fitted radius model

$$\frac{R}{R_E} = A \left[ \log_{10} \left( \frac{\mathcal{L}}{\mathcal{L}_{\odot}} \right) \right]^4 + B \left[ \log_{10} \left( \frac{\mathcal{L}}{\mathcal{L}_{\odot}} \right) \right]^3$$

$$+ C \left[ \log_{10} \left( \frac{\mathcal{L}}{\mathcal{L}_{\odot}} \right) \right]^2 + D \left[ \log_{10} \left( \frac{\mathcal{L}}{\mathcal{L}_{\odot}} \right) \right] + E$$

- Inflation response is instantaneous

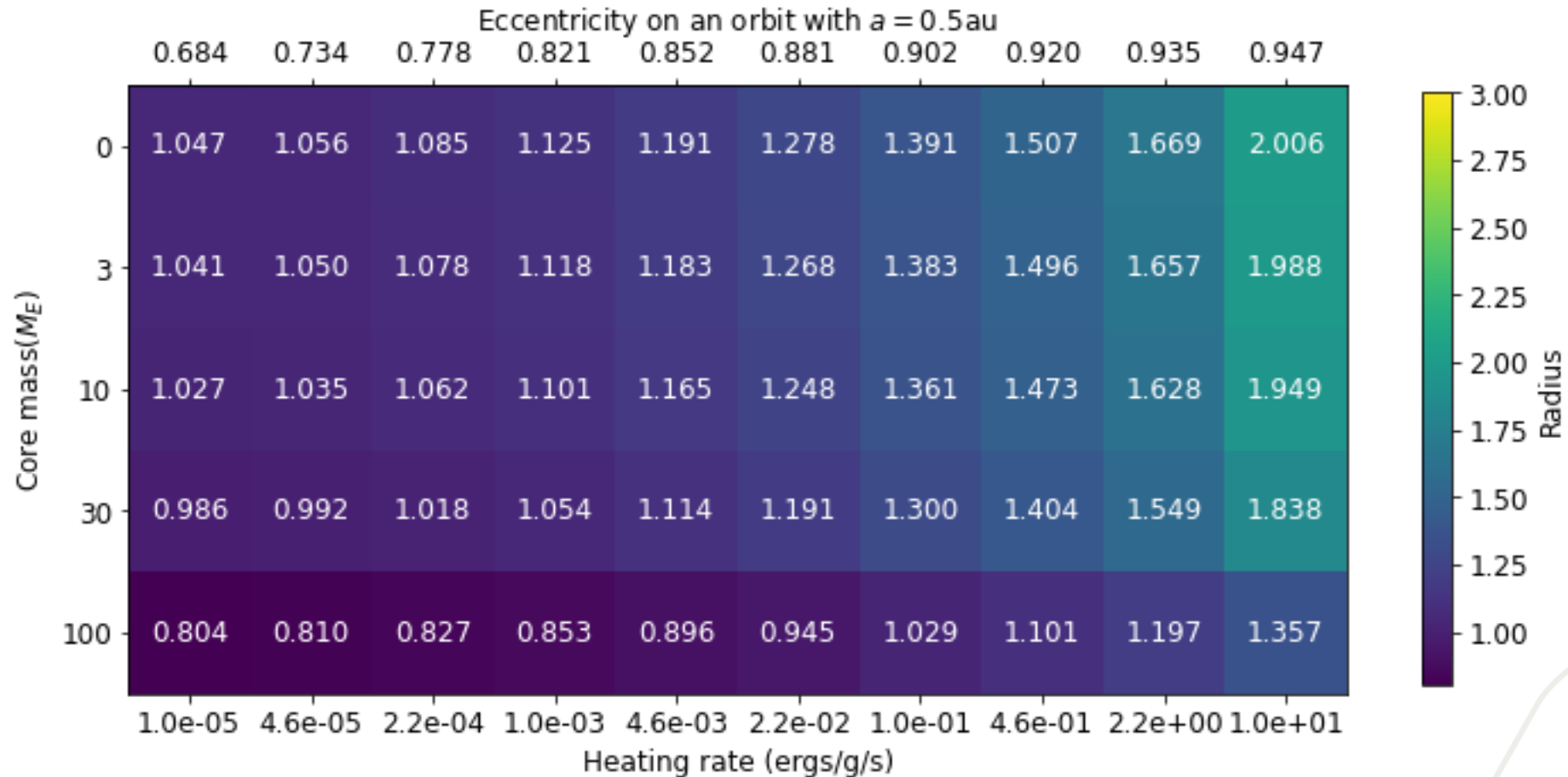




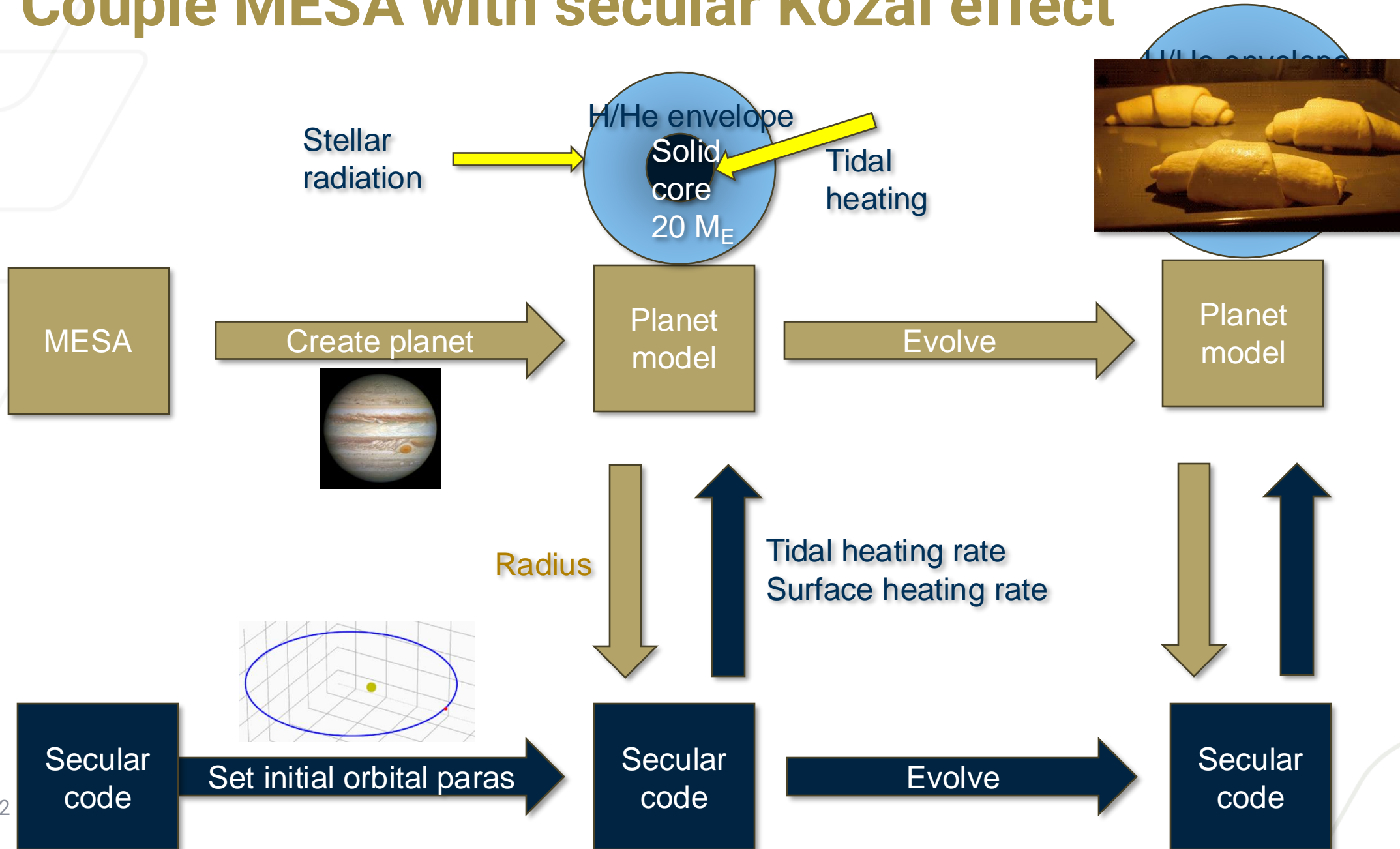
# Dependence on core mass

- Constant heating on a Jupiter mass planet
  1. Smaller core mass ---> larger radius
  2.  $a=0.5$  au,  $e=0.95$  --->  $2 R_J$

Note: We use  $20 M_E$  core mass for our case

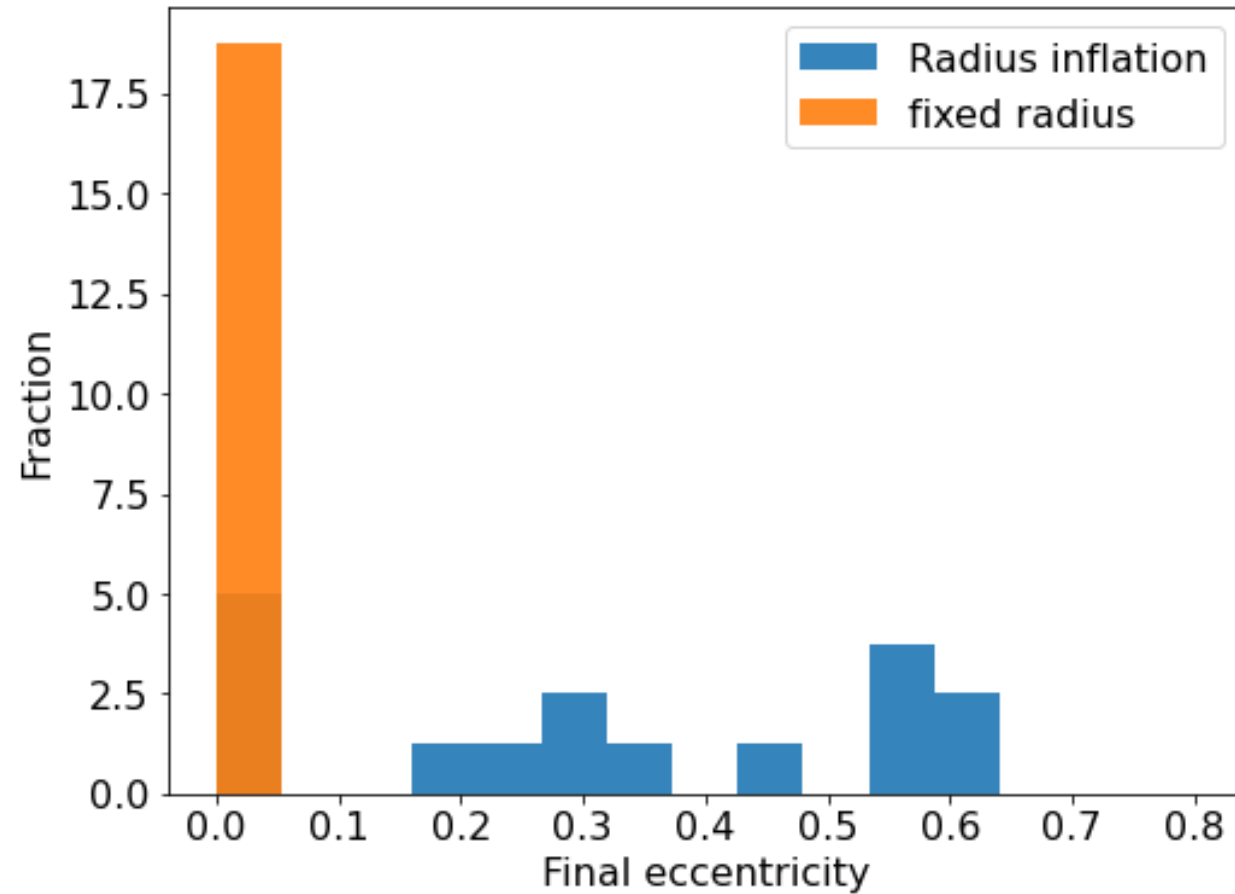


# Couple MESA with secular Kozai effect

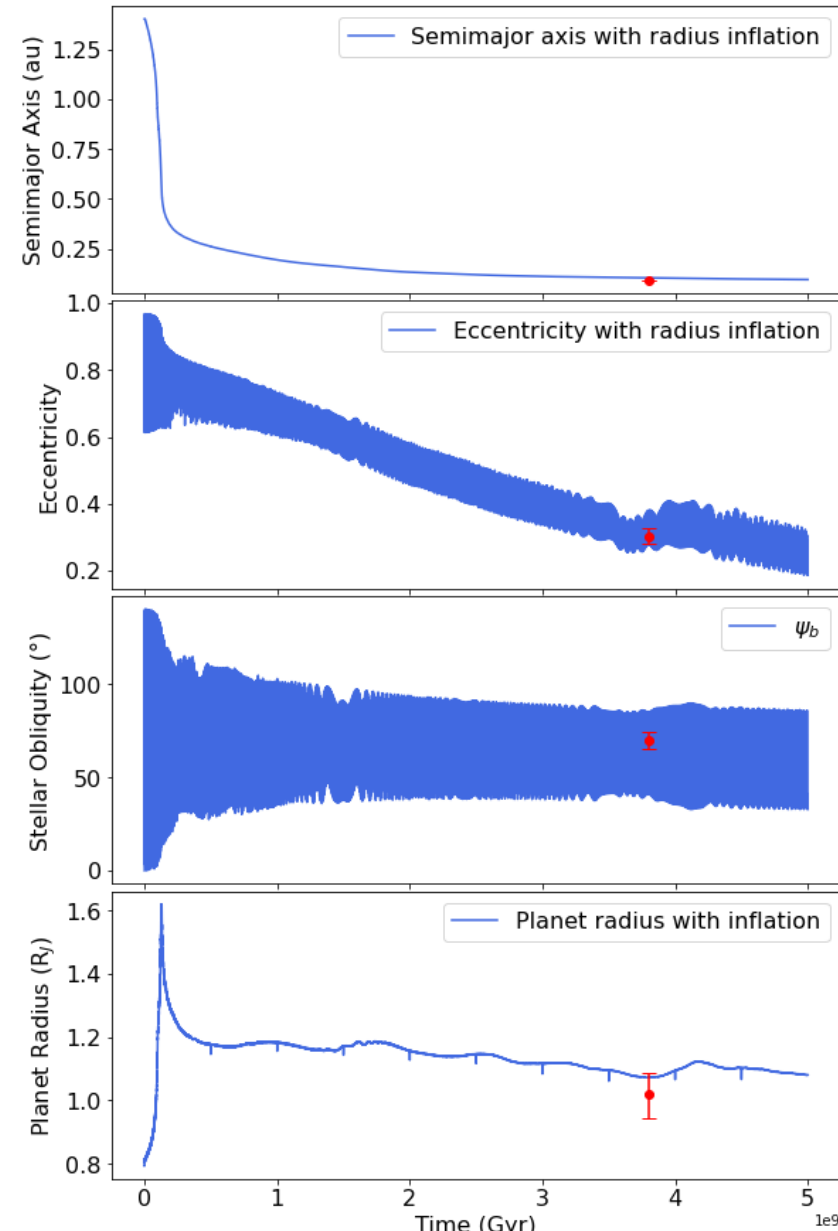
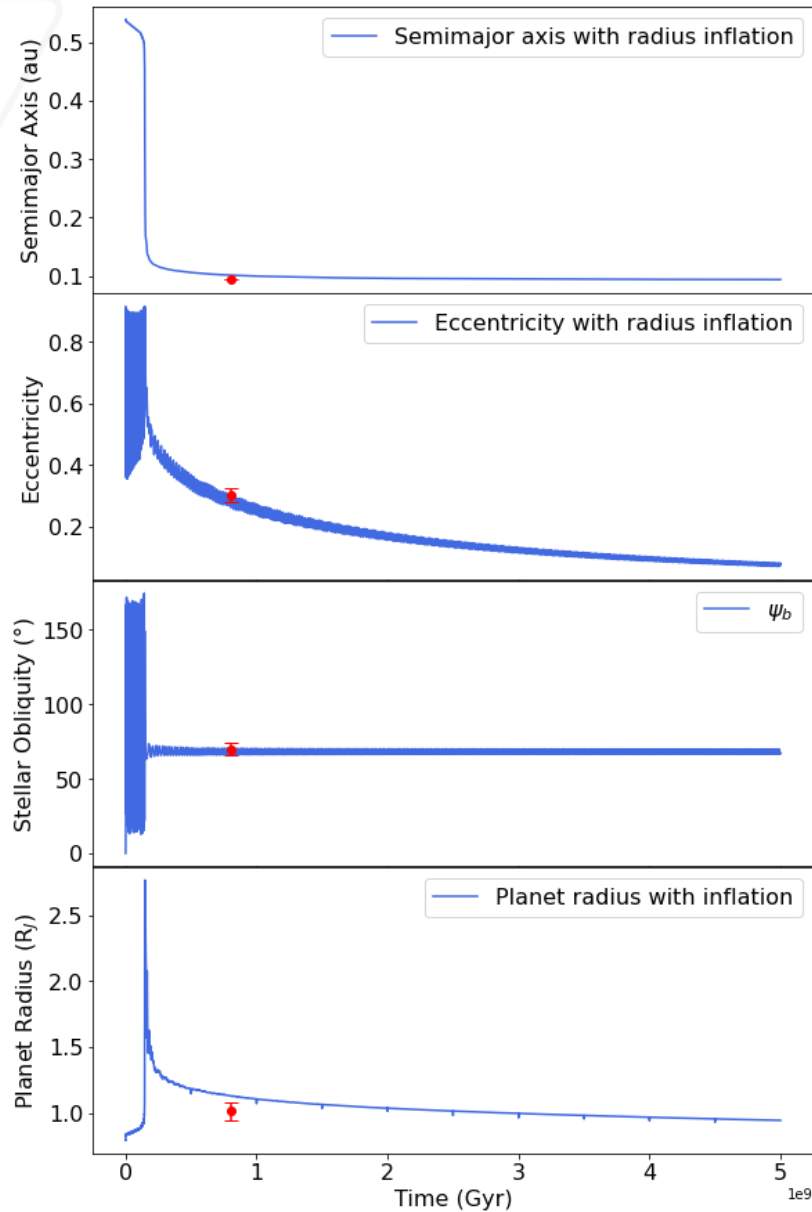


# 5 $M_E$ core mass simulation

- Final eccentricity

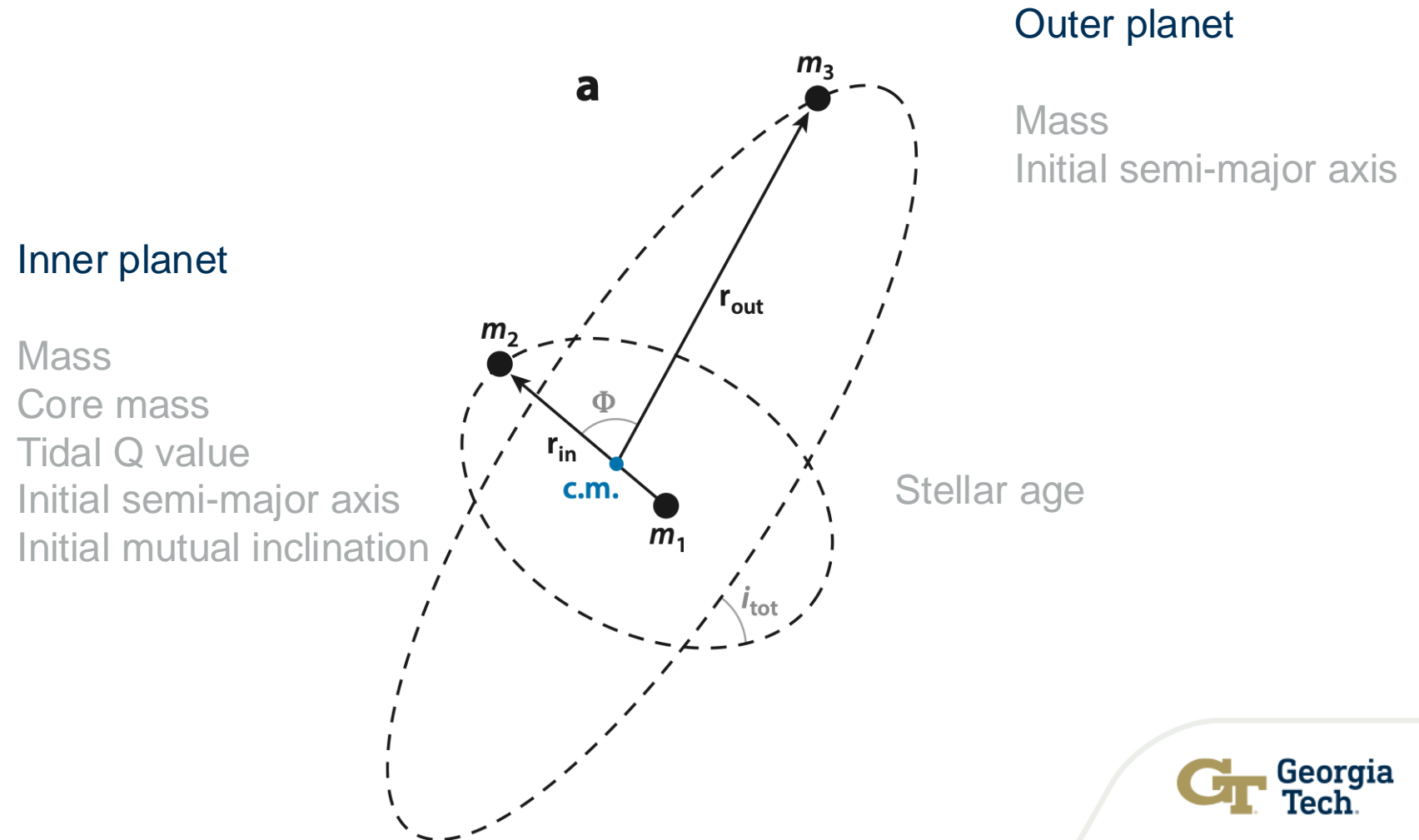


# More examples

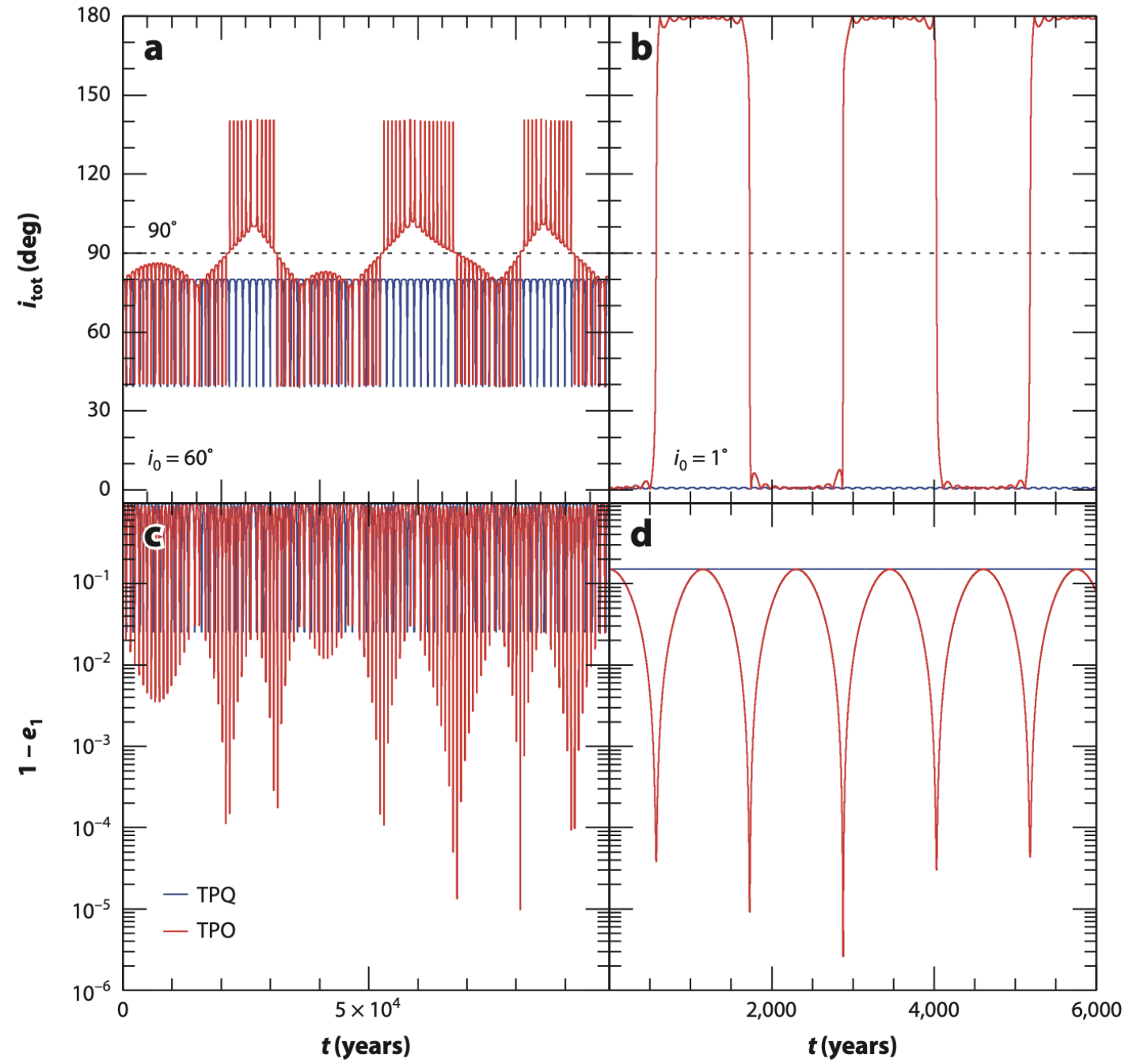
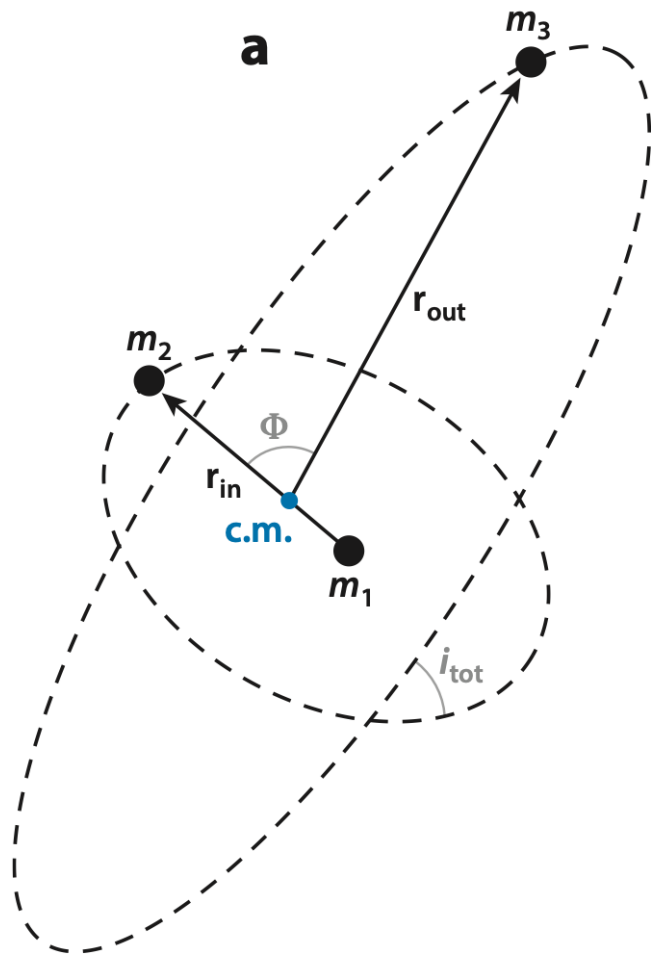


# Constrain parameters of a specific system (Such as WASP 117b)

- Large parameter space



# Octupole L-K



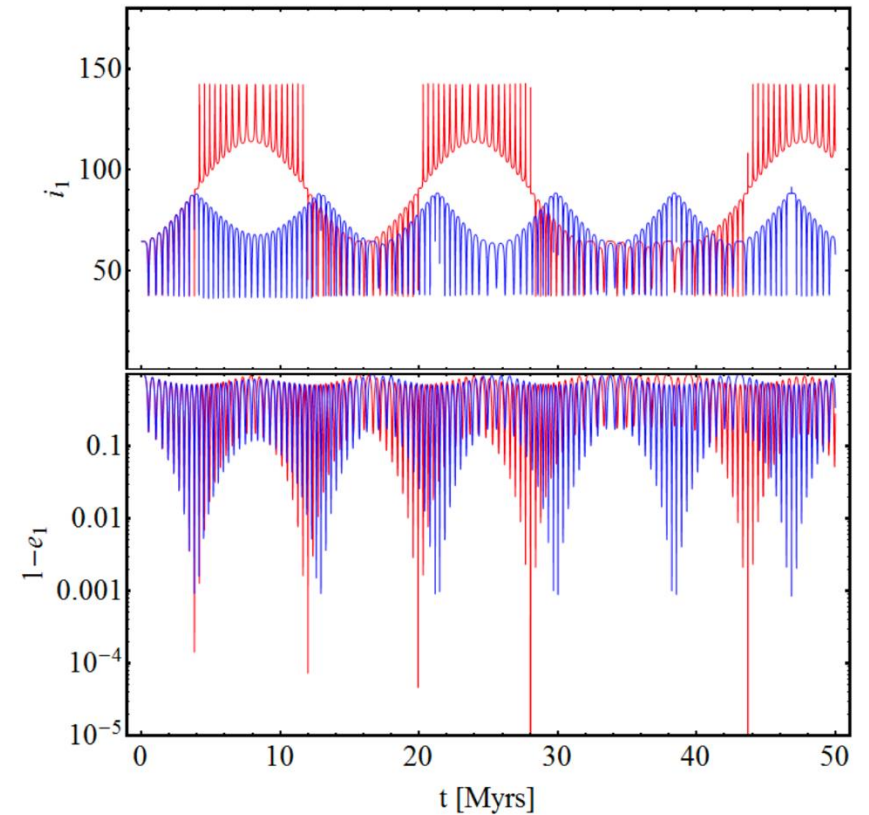


# Short range forces

- Stop planet from plunge into the star

$$\varepsilon_{\text{GR}} \equiv \frac{3G(m_0 + m_1)^2 a_2^3 (1 - e_2^2)^{3/2}}{a_1^4 c^2 m_2}$$

$$\varepsilon_{\text{Tide}} \equiv \frac{15m_0(m_0 + m_1)a_2^3(1 - e_2^2)^{3/2}k_{2,1}R_1^5}{a_1^8 m_1 m_2}$$



# Tidal dissipation

- Tidal dissipation rate:

$$L_{\text{tide}}(e, \epsilon) = \frac{2K}{1 + \cos^2 \epsilon} [\sin^2 \epsilon + e^2 (7 + 16 \sin^2 \epsilon)].$$

$$K = \frac{3n k_2}{2 Q} \left( \frac{GM_{\star}^2}{R_p} \right) \left( \frac{R_p}{a} \right)^6,$$

- High eccentricity will introduce super high heating rate, make the planet explode.